## 26. Molecular Ions of Transient Species: Vinyl-Alcohol Cation

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## Summary

Vinyl alcohol 1 was prepared by thermolysis of cyclobutanol and its photoelectron spectrum was determined.  $I_a^1 = 9.18$  eV and  $I_v^1 = 9.52$  eV were found, the vibrational progression ( $\tilde{\nu} = 1400$  cm<sup>-1</sup>) for this lowest energy transition  $1(X) \rightarrow 1^+(\tilde{X})$  indicating significant skeletal changes in the ion. The question of the relative stability of the *syn*  $(I_a^+)$ -*vs. anti*-ions  $(I_a^+)$  is discussed in the light of theoretical calculations. The energy of the second  $\pi$ -state of  $1^+$  is estimated at 13.6–14.1 eV above the ground state of 1.

Introduction. – Vinyl alcohol 1, the unstable tautomer of acetyldehyde 2, has been suggested for a century [1] as an intermediate in various chemical reactions. Only ten years ago, it was identified for the first time in solution [2], and subsequently in the gas phase [3] where its molecular structure could be derived from its microwave spectrum. The existence of related species in interstellar space (*e.g.* 2, CH<sub>3</sub>OH, C<sub>2</sub>H<sub>5</sub>OH, CH<sub>2</sub>CHCN) suggests, however, that 1 may also be detected there provided sufficient knowledge about its molecular constants is available [3]. Detection of its radical cation 1<sup>+</sup> may even be more likely given that its heat of formation is 62.7 kJmol<sup>-1</sup> lower than the one of isomeric 2<sup>+</sup> [4], contrary to the situation prevailing for the neutrals. In this paper the electronic states of 1<sup>+</sup> are investigated using photoelectron (PE) spectroscopy.

**Experimental.** - The details of the PE setup have been described in [5]. 1 was prepared by pyrolysis of cyclobutanol 3 [4] in a Pt/Rh-tube ( $\emptyset = 0.9$  mm) containing a silica capillary insert and attached directly to the target chamber. This silica insert proved to be crucial since pyrolysis in the hot metal tube alone led exclusively to 2 and ethylene. Heating was achieved by applying a controlled voltage at the tube (used as a resistor); 550-600° were found to be optimal for decomposition of 3. The temperature of the fragments (cooled down through collisions with the walls) in the ionization chamber was estimated to be 100°-150° from the intensity of the hot vibrational band in the third PE transition of naphthalene (= 'internal molecular thermometer'), investigated under the same conditions and compared to the results given in [5]. Calibration of  $\approx 30$  meV was determined.



Figure. a) PE spectrum of pyrolysed cyclobutanol (see Experimental). b) Spectrum a from which the PE spectra of ethylene and acetaldehyde were digitally subtracted, employing weight factors calibrated from the intensity of the first PE peaks of these by-products from a. The spectrum corresponds essentially to that of vinyl alcohol, apart from spurious peaks due to a non-perfect subtraction. c) Band system ① of vinyl alcohol, expanded scale.

**Results and Discussion.** – The *Figure* shows the PE spectrum of 1 together with that of ethylene and isomeric 2, the latter being apparently formed through catalytic action of the oven surface. 1 possesses two filled  $\pi$ -orbitals; the subsequent discussion is restricted to the two ionic <sup>2</sup>A"( $\pi$ )-states of 1<sup>+</sup> derived from ejection of an electron out of these two orbitals.

Band System  $\bigcirc$ .  $I_v^{l} = 9.52$  eV can be deduced from band system  $\bigcirc$ . This value is significantly higher than those determined for alkylvinylethers CH<sub>2</sub>CHOR ( $I_v^{l} = 9.14$ eV (R = Me); = 9.15 eV (R = Et); = 9.10 eV (R = Bu) [6]). This finding is in line with the fact that from perturbation theory based on iso- $\pi$ -electronic allyl anion as well as from MNDO calculations, the HOMO of 1 is strongly localized in the C = C-fragment. Various alkyl substituents at the O-atom thus provide a rather constant inductive destabilization to this orbital ( $\approx 0.4 \text{ eV}$ ) since they do not significantly engage in ionic hyperconjugation. The MNDO-prediction  $I_v^{l} = 9.4$  eV (applying *Koopmans'* theorem) is in good agreement with experiment. This value was calculated for the minimum energy structure of 1 which – in line with experiment [3] – turned out to be the *syn*-form ( $\mathbf{1}_{s}$ ).

Band <sup>(1)</sup> displays a long vibrational progression of 1400 cm<sup>-1</sup> indicating that considderable nuclear rearrangement takes place upon  $1 \rightarrow 1^+ + e^-$ . Qualitative arguments as well as numerical calculations [7] predict for this process a lengthening of the C=C and a shortening of the C-O bonds (from [7]:  $\Delta R_{C=C} = +10$  pm,  $\Delta R_{C=O} = -10$  pm). The observed frequency is not unexpected for a corresponding fundamental excited in the ion, also in view of the C=C stretching frequency in iso- $\pi$ -electronic allyl radical (1477 cm<sup>-1</sup> [8]).

A recent investigation of the appearence potential of 1<sup>+</sup> using mass-spectrometric techniques has shown a low cross section for ionization over the initial portion of the ion production efficiency curve [4]. Small Franck-Condon (FC) factors for an adiabatic process  $1 \rightarrow 1^+$  were made responsible for this observation, and rationalized through the theoretical prediction [7] that – contrary to the most stable conformation for 1 (*i.e.* 1,) - the ion 1<sup>+</sup> prefers the *anti*-conformation (*i.e.*  $1_a^+$ ). Indeed MNDO yields 10.9 kJmol<sup>-1</sup> for the difference  $\Delta H_{f}^{o}(\mathbf{1}_{s}^{+}) - \Delta H_{f}^{o}(\mathbf{1}_{a}^{+})$ . However, MNDO also yields a significant barrier for  $\mathbf{1}_{s}^{+} \rightarrow \mathbf{1}_{s}^{+}$  of 96.5 kJmol<sup>-1</sup>, in contrast to the corresponding value for the neutrals of  $10.5 \text{ kJmol}^{-1}$ . Hence, the ion potential wells are well-developed with respect to this torsional coordinate. In fact one might ask if  $1_s^+$  and  $1_a^+$  can be addressed as 'rotamers' in the usual sense (i.e. low barriers to interconversion), or if the notion 'isomers' would not better apply. We note in this context that the resonance energy of iso- $\pi$ -electronic vinylamine cations is estimated experimentally to be  $\approx 1 \text{ eV}$  [9]. The PE transition  $\bigcirc$ therefore most likely corresponds to  $\mathbf{1}_s \rightarrow \mathbf{1}_s^+$ . Its vibrational activity due to the already discussed in-plane skeletal change indicates a considerable non-adiabaticity which might be the reason for the low ion production efficiency noted in [4].

In our view neither of the two techniques (guided by the FC principle) has yielded information about the energy of the most stable form  $1_a^+$ , *i.e.* about  $I_a^1$   $(1_s \rightarrow 1_a^+)$ . A value for  $I_a^1$   $(1_s \rightarrow 1_s^+)$ , can, however, be derived from the shoulder at 9.18 eV, in good agreement with [4]. This shoulder has been reproduced several times and its intensity did not vary with decomposition temperature. Furthermore it appears as a member of the 1400 cm<sup>-1</sup> progression. These arguments exclude it from being a hot band, given also the effective temperature of 1 (see *Experimental*). Finally, a signal due to a small thermal equilibrium population of  $1_a$  can be excluded as the microwave spectrum obtained under similar conditions [3] did not reveal the presence of this rotamer.

Band System @. MNDO predicts two  ${}^{2}A'(\sigma)$ -states ( $\tilde{A}$  and  $\tilde{B}$ ) of 1<sup>+</sup> at 13.42 and 13.57 eV, and a second  ${}^{2}A''(\pi)$ -state ( $\tilde{C}$ ) at 14.12 eV. Another crude estimate is obtained through the PE-LCBO model [10], using  $H_{C=C} = -10.5$  eV (from ethylene) and  $H_{\tilde{O}} = -12.6$  eV (from water). With the experimental value  $\varepsilon_{HOMO}$  (1) = -9.52 eV,  $\varepsilon = -13.58$  eV for the subjacent  $\pi$ -MO is derived.

There are ionic states present in the energy region 13.6–14.1 eV, but a clear-cut assignment of the second  ${}^{2}A''(\pi)$ -state of 1<sup>+</sup> is not possible due to overlapping bands from the  $\sigma$ -states, as well as non-perfect subtraction of the bands from the pyrolysis by-products. Nevertheless, as a rough estimate for the strongly allowed  $\tilde{C} \rightarrow \tilde{X}$  emission of the ion 1<sup>+</sup> the UV region 250–300 nm looks acceptable. However, from the lower lying excited  $\sigma$ -states  $\tilde{A}$  and  $\tilde{B}$  emission is also weakly allowed for this C<sub>s</sub>-species.

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## REFERENCES

- [1] E. Erlenmeyer, Chem. Ber. 14, 320 (1881).
- [2] B. Blank & H. Fischer, Helv. Chim. Acta 56, 506 (1973).
- [3] S. Saito, Chem. Phys. Lett. 42, 399 (1976).
- [4] J. L. Holmes & F. P. Lossing, J. Am. Chem. Soc. 104, 2648 (1982); see also R. A. Back, Can. J. Chem. 60, 2537 (1982).
- [5] R. Dressler, L. Neuhaus & M. Allan, J. Electron Spectrosc. 31, 181 (1983).
- [6] H. Friege & M. Klessinger, J. Chem. Res. (S) 1977, 208.
- [7] W.J. Bouma, J.K. Mac Leod & L. Radom, J. Am. Chem. Soc. 101, 5540 (1979).
- [8] G. Maier, H. P. Reisenauer, B. Rohde & K. Dehnicke, Chem. Ber. 116, 732 (1983).
- [9] R. Houriet, J. Vogt & E. Haselbach, Chimia 34, 277 (1980).
- [10] H. Bock, G. Wagner, K. Wittel, J. Sauer & D. Seebach, Chem. Ber. 107, 1869 (1974); see also H. Bock & B. Solouki, Angew. Chem. 20, 425 (1981); H. Bock, Chem. Rundschau 29, 208 (1981).